

1. 135(3). Antisymmetry laws without the constraint $\underline{A} = \underline{0}$.

∂_n & $u(1)$ level:

$$\underline{E} = -\underline{\nabla} \phi - \frac{\partial \underline{A}}{\partial t}, \quad - (1)$$

$$\underline{B} = \underline{\nabla} \times \underline{A}. \quad - (2)$$

They are now constrained by:

$$\underline{\nabla} \phi = \frac{\partial \underline{A}}{\partial t}, \quad - (3)$$

$$\frac{\partial A_i}{\partial x_j} = - \frac{\partial A_j}{\partial x_i}, \quad - (4)$$

It follows that:

$$\underline{\nabla} \times \underline{\nabla} \phi = \underline{\nabla} \times \frac{\partial \underline{A}}{\partial t} = \underline{0}, \quad - (5)$$

so:

$$\underline{B} = \underline{\nabla} \times \underline{A} = \underline{0}. \quad - (6)$$

The antisymmetry constraints (3) and (4) cause the magnetic field to vanish. So $u(1)$ electrodynamics is restricted by antisymmetry to electric fields only. In this note a potential is derived which is associated with the antisymmetry rule. The only situation in which the electric field does not vanish is defined by this potential:

$$\oint \underline{A} \cdot d\underline{\ell} = \int_S (\underline{\nabla} \times \underline{A}) \cdot \underline{n} dA \quad - (7)$$

$$= 0,$$

i.e. by

$$\underline{\nabla} \times \underline{A} = \underline{0}, \quad - (8)$$

and by:

$$\frac{\partial A_x}{\partial y} = -\frac{\partial A_y}{\partial x} \quad \text{etc.} \quad - (9)$$

Resolvo: $\underline{A} = A_x \underline{i} + A_y \underline{j} + A_z \underline{k} \quad - (10)$

It also obeys the antisymmetry equation:
 $\frac{\partial \phi}{\partial z} = \frac{\partial A_z}{\partial t} \quad \text{etc.} \quad - (11)$

Travelling Waves

Attempt the solution:
 $\phi = \frac{\phi^{(0)}}{\sqrt{2}} \exp(i(\omega t - \kappa z)) \quad - (12)$

then: $\frac{\partial \phi}{\partial z} = \frac{\partial A_z}{\partial t} = -i\kappa \phi \quad - (13)$

$$A_z = -i\kappa \int \phi dt \quad - (14)$$

In general: $\underline{\nabla} \phi = -i\kappa \phi \left(\underline{i} + \underline{j} + \underline{k} \right) \quad - (15)$

$$\frac{\partial \underline{A}}{\partial t} = -i\frac{\omega}{c} \phi \left(\underline{i} + \underline{j} + \underline{k} \right) \quad - (16)$$

$$i(\omega t - \kappa z) \quad - (17)$$

$$\underline{E} = 2i\kappa \phi \left(\underline{i} + \underline{j} + \underline{k} \right) e$$

$$\text{Re}(\underline{E}) = -2\frac{\kappa}{\omega} \phi \sin(\omega t - \kappa z) \left(\underline{i} + \underline{j} + \underline{k} \right) \quad - (18)$$

In free space on the u(1) level:

$$\underline{\nabla} \times \underline{E} + \frac{\partial \underline{B}}{\partial t} = \underline{0} \quad - (19)$$

$$\underline{\nabla} \times \underline{B} - \frac{1}{c^2} \frac{\partial \underline{E}}{\partial t} = \underline{0} \quad - (20)$$

$$\underline{\nabla} \cdot \underline{B} = 0 \quad - (21)$$

$$\underline{\nabla} \cdot \underline{E} = 0 \quad - (22)$$

A solution of the type (18) does not say:

$$\underline{\nabla} \times \underline{E} = 0, \quad \frac{\partial \underline{E}}{\partial t} = 0, \quad \underline{\nabla} \cdot \underline{E} = 0 \quad - (23)$$

so cannot be a free space solution. If for example:

$$\underline{E} = E_2 \underline{k}, \quad - (24)$$

$$E_2 = 2 \kappa \omega \phi e^{i(\omega t - \kappa z)} \quad - (25)$$

$$\begin{aligned} \underline{\nabla} \cdot \underline{E} &= \frac{\partial E_2}{\partial z} = 2 \kappa^2 e^{i(\omega t - \kappa z)} \phi \\ &= \rho / \epsilon_0 \quad - (26) \end{aligned}$$

Therefore:

$$\rho = 2 \epsilon_0 \kappa^2 \phi e^{i(\omega t - \kappa z)} \quad - (27)$$

Similarly:

$$-\frac{1}{c^2} \frac{\partial \underline{E}}{\partial t} = \mu_0 \underline{J} \quad - (28)$$

From eq. (25):

$$\frac{\partial E_2}{\partial t} = -2 \kappa \omega \phi e^{i(\omega t - \kappa z)} \quad - (29)$$

So:

$$\underline{J}_z = \frac{2 \kappa \omega \phi}{\mu_0 c^2} e^{i(\omega t - \kappa z)} \quad - (30)$$

4) If this wave propagates at a velocity v :

$$k = \frac{\omega}{v} \quad - (31)$$

so:

$$p = 2\epsilon_0 \frac{\omega^2}{v^2} \phi e^{i(\omega t - kz)} \quad - (32)$$

$$J_z = 2\epsilon_0 \frac{\omega^2}{v} \phi e^{i(\omega t - kz)} \quad - (33)$$

or average:

$$\langle p \rangle = 0 \quad - (34)$$

$$\langle J_z \rangle = 0 \quad - (35)$$

and:

$$\langle E \rangle = 0 \quad - (36)$$

The root mean squares are:

$$\langle p^2 \rangle^{1/2} = \epsilon_0 \left(\frac{\omega}{v}\right)^2 \phi \quad - (37)$$

$$\langle J_z^2 \rangle^{1/2} = \epsilon_0 \frac{\omega^2}{v} \phi \quad - (38)$$

$$\langle E_z^2 \rangle^{1/2} = k \phi \quad - (39)$$

$$= \frac{\omega}{v} \phi \quad - (40)$$

Coulomb Law

In this case:

$$\phi = - \frac{e}{4\pi\epsilon_0 r} \quad - (41)$$

$$\frac{\partial \phi}{\partial r} = \frac{\partial A_z}{\partial t} = \frac{e}{4\pi\epsilon_0 r^2} \quad - (42)$$

>), So:

$$A_z = \int \frac{e}{4\pi\epsilon_0 z^2} dt \quad - (43)$$

i.e $A_z(t) = \frac{et}{4\pi\epsilon_0 z^2} + A_z(0) \quad - (44)$

$$A_z(t) - A_z(0) = \frac{et}{4\pi\epsilon_0 z^2} \quad - (45)$$

Define: $v = \frac{z}{t} \quad - (46)$

then: $\phi = -vA_z \quad - (47)$

$$E_z = -\frac{2e}{4\pi\epsilon_0 z^2} \quad - (48)$$

ECE Level if $A \neq 0$

The result is:

$$\underline{E}^a = \underline{E}_A^a + \phi^{(0)} \frac{\omega}{E} \quad - (49)$$
$$\underline{B}^a = \frac{\phi^{(0)}}{c} \frac{\omega}{B} \quad - (50)$$

where \underline{E}_A^a is due to a non-zero A.